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Effect of Ta seed layer on crystalline structure and magnetic properties in an exchange-biased Co/IrMn system

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ABSTRACT

Ta-seeded and un-seeded layers of a top-configuration Co/IrMn system were deposited onto glass substrate by DC sputtering. Three sets of deposition conditions for $Co(50 \text{ Å})/IrMn(t_{IrMn} \text{ Å})$ and $Co(t_{Co} \text{ Å})/IrMn(90 \text{ Å})$, where $t_{IrMn} = 15$, 30, 60, 90, 110, and 150 Å, and $t_{Co} = 15$, 25, 50, 75, 100, 125, and 150 Å, were: condition (a) substrate temperature (T_s) was kept at room temperature (RT). Condition (b) T_s set to RT, with in-plane magnetic field, $H = 500 \, \text{Oe}$. In condition (c), condition (b) was followed by post-deposition annealing in the magnetic field at $T_A = 250 \,^{\circ}\text{C}$ for 1 h, then field cooled to RT. X-ray diffraction (XRD) patterns and grazing incidence scans revealed maximum IrMn (111) texture to occur for post-deposition annealed Ta seed layer samples. The IrMn (111) texture-effect significantly influences magnetic properties, including exchange-biasing field (H_{ex}) , interfacial energy (J_k) , and coercivity (H_c) . The Ta seed layer also significantly influences magnetic properties. Adding a Ta seed layer to the Co/IrMn system increases Hex, because of the IrMn (1 1 1) texture. For Ta-seeded Co/IrMn under condition (c), $H_{\rm ex}$ tended to saturate for $t_{\rm IrMn} \ge 90$ Å. Under conditions (a) and (b), $H_{\rm ex}$ decreased with increasing t_{IrMn} for $t_{\text{IrMn}} \ge 90 \,\text{Å}$. H_{ex} values for all un-seeded Co/IrMn systems increased with t_{IrMn} . J_k versus t_{IrMn} plot is proportional only to H_{ex} in the Ta-seeded and un-seeded layers of a top-configuration Co/IrMn system, due to the interfacial energy formula, t_{Co} is fixed, and saturation magnetization (M_s) of the Co layer is constant. Results for the Ta-seeded system showed a strong relationship between H_c and t_{IrMn} , due to coupling-decoupling interactions between Co spin, and IrMn layers close to the Co/IrMn interface. The $H_{\rm ex}$ versus $t_{\rm Co}$ result shows that the $H_{\rm ex}$ is proportional to $(1/t_{\rm Co})$. The $H_{\rm ex}$ values with the Ta seed layer are almost slightly larger than those without a Ta seed-layer. The dependence of J_k on t_{Co} is similar to the trend in M_s on t_{Co} , I_k tends to saturate slowly as t_{Co} increases. Surface pinning occurred in all systems, revealing an inverse relation between H_c and t_{Co} . Removing the Ta seed-layer weakens IrMn (111) texturing, reducing $H_{\rm ex}$. The maximum observed $H_{\rm ex}$ and $J_{\rm k}$ values were 205 Oe and 0.11 erg/cm², respectively.

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1. Introduction

Magnetoresistance random access memory (MRAM) and recording-head applications could exploit ferromagnetic (FM)/antiferromagnetic (AFM) interface exchange-coupling interactions, which cause a unidirectional shift in the magnetic hysteresis loop exchange-bias [1–5]. Numerous factors, including the IrMn (111) texture, the AFM layer ordered and disordered phases, post-annealing scheme, and the seed layer effect [6–9], can influence FM/AFM magnetic properties. Researchers attributed

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the larger exchange-biasing field ($H_{\rm ex}$) to a stronger IrMn (111) texture, where the seed layer deposition-candidate is crucial [10–12]. There is a strong relationship between $H_{\rm ex}$ performance and IrMn layer characteristics, including (111) texture, anisotropy constant ($K_{\rm AF}$), grain size distribution, and AFM grain thermal-stability [10–14]. The seed layer is probably the dominant factor affecting microstructural orientation, reduction in current shunting for grain magnetoresistance (GMR), and the tunneling magnetoresistance (TMR) effect [13,14]. Moreover, removing the Ta seed layer from the Co/IrMn system results in poor IrMn (111) texture, and reduces $H_{\rm ex}$ values. Thus, there is a close relationship between the seed layer effect, and magnetic performance variables such as exchange-biasing system $H_{\rm ex}$ and coercivity ($H_{\rm c}$).

We deposited a top-configuration Co/IrMn system onto glass substrate under three different deposition conditions, described in the following experimental procedure section. This investigation

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focuses on the relationship between IrMn (111) texture and the materials magnetic properties.

2. Experimental procedures

The top configuration Co/IrMn system was fabricated on glass substrate by DC magnetron sputtering. We considered two cases: (A) devices without a Ta seed layer and (B) devices with a Ta seed layer. The Ta-seeded and un-seeded multilayer deposition sequences were, respectively: (1) Glass/Ta(30 Å)/Co(50 Å)/IrMn(t_{IrMn} Å)/Ta(100 Å) and Glass/ $Co(50 \text{ Å})/IrMn(t_{IrMn} \text{ Å})/Ta(100 \text{ Å})$, where t_{IrMn} = 15, 30, 60, 90, 110, and 150 Å. (2) $Glass/Ta(30\,\text{\AA})/Co(t_{Co}\,\text{Å})/IrMn(90\,\text{Å})/Ta(100\,\text{Å}) \quad \text{ and } \quad Glass/Co(t_{Co}\,\text{Å})/IrMn(90\,\text{Å})/Ta(100\,\text{Å})$ Ta(100 Å), where $t_{Co} = 15$, 25, 50, 75, 100, 125, and 150 Å. The top Ta(100 Å) layer served as a protective layer to prevent oxidation of IrMn. For each system, we applied three different sets of deposition conditions: (a) the substrate temperature (T_s) was maintained at room temperature (RT) in the absence of any magnetic field. (b) T_s was RT, with an in-plane magnetic field, $H = 500 \,\mathrm{Oe}$, during deposition. (c) Condition (b) was followed by post-deposition annealing in the field at 250 °C (T_A) for 1 h, then cooled to RT in the magnetic field. The use of Ta as a seed layer produces a strong (111) texture in the Co and IrMn layers [15]. The composition of the IrMn alloy target was 20 at.% Ir and 80 at.% Mn. Typical base chamber pressure was 2×10^{-7} Torr, and the argon gas working pressure was 5×10^{-3} Torr.

The degree of $Ir_{20}Mn_{80}$ (111) layer texturing was characterized by X-ray diffraction (XRD) and grazing incidence (GIXD) scans at a grazing angle of 1°, using Cu $K\alpha_1$ radiation. The LakeShore Model 7300 vibrating sample magnetometer (VSM) provided the exchange-biased magnetic hysteresis loop for training effect removal.

3. Results and discussion

First, we emphasized the two cases (A and B) again, respectively. The case A means that the Co/IrMn system is no Ta seed layer structure. In contrast, the case B means that the Co/IrMn system has the Ta seed layer structure. Fig. 1(a) shows the Co(50 Å)/IrMn(110 Å) XRD patterns, for each case (A and B) in the same post-annealing process. Fig. 1(a) compares the X-ray results for the two cases. The Co/IrMn system with a Ta seed layer (case B), has a crystalline bodycentered cubic (BCC) structure, producing β -Ta (200) and Ta (110) peaks. Moreover, the diffraction intensities of Co (111) and IrMn (111) planes in case B were much stronger than those in case A, indicating that the inclusion of a Ta seed layer, results in the Co and IrMn layers having more crystalline character. Fig. 1(b) shows the corresponding grazing angle scans to Fig. 1(a). The spectrum shows that the annealed Ta seed-layer sample produces strong IrMn (111) and Co(111) in-plane diffraction-peaks. By contrast, in the absence of a Ta seed layer in the annealed sample, only weak IrMn (111) and Co (111) peaks appear. Fig. 1(c) shows that IrMn (111) X-ray peak intensity is dependent on the thickness (t_{IrMn}) of IrMn in each case. Post-annealing treatment of case B (with a Ta seed layer) samples produces the strongest IrMn (111) texture, while the IrMn (111) peaks of the remaining samples are weak and less intense. Accordingly, adding a Ta seed layer, without post-annealing treatment, does not necessarily strengthen IrMn texturing.

Fig. 2(a) is a plot of the exchange field $(H_{\rm ex})$, as a function of IrMn thickness $(t_{\rm IrMn})$. This plot shows that the Ta seed layer significantly influences magnetic properties. Adding a Ta seed layer to the Co/IrMn system increases $H_{\rm ex}$ due to the IrMn (1 1 1) plane, that provides the largest possible spin density at Co/IrMn interfaces [16]. Therefore, $H_{\rm ex}$ in case A (without a Ta seed layer) is generally lower than that in case B (with a Ta seed layer). In both cases A and B, $H_{\rm ex}$ increases with $t_{\rm IrMn}$, from 15 Å to 60 Å. $H_{\rm ex}$, for Ta-seeded Co/IrMn under condition (c), tends to saturate at $t_{\rm IrMn} \ge 90$ Å, while for Ta-seeded Co/IrMn under conditions (a) and (b), $H_{\rm ex}$ tends to decline as $t_{\rm IrMn}$ increases over this range, because of the IrMn texture defect [11]. By contrast, the $H_{\rm ex}$ value for all the un-seeded Co/IrMn systems tends to increase slightly with $t_{\rm IrMn}$. Fig. 2(b) plots interfacial energy $(J_{\rm k})$ versus $t_{\rm IrMn}$, using a well-established formula,

$$J_{k} = H_{ex} M_{s} t_{Co}, \tag{1}$$

where M_S is the saturation magnetization of the Co layer and t_{Co} is the Co layer thickness. Since t_{Co} is fixed 50 Å in both cases A and B,

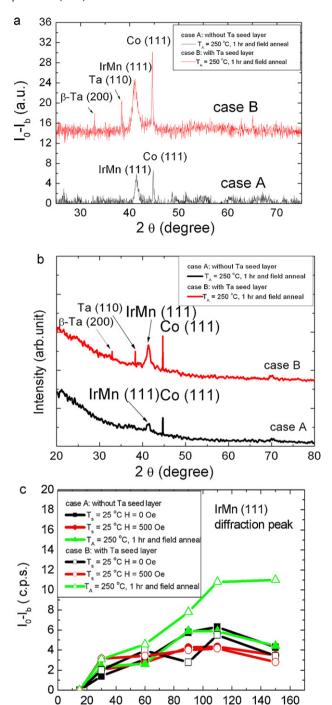


Fig. 1. (a) X-ray diffraction plots in cases A and B. I_o is the absolute intensity and I_b is the background intensity. (b) Grazing angle result for post-annealing samples with cases A and B. (c) Degree of IrMn (111) texture, as determined from the X-ray diffraction results, plotted as function of $t_{\rm IrMn}$ for cases A and B. I_o is the intensity of the IrMn (111) line and I_b is the background intensity.

t_{IrMn} thickness (Å)

and $M_{\rm S}$ is constant, then, $J_{\rm k}$ is proportional only to $H_{\rm ex}$, as shown in Fig. 2(b). According to the above results of the XRD analysis and exchange-biasing magnetic hysteresis loop measurement, not only the density of spins at the Co/IrMn interface, but also $J_{\rm k}$ increases due to the presence of Ta seed layer, and to the field annealing treatment. This indicates that the effects of both Ta seed layer, and field annealing treatment, on exchange-coupling interactions at the Co/IrMn interface are significant and effective. In view of literature reports regarding exchange-biasing, higher $J_{\rm k}$ can over-

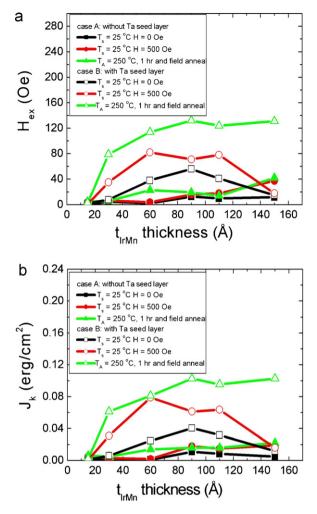


Fig. 2. (a) Dependence of the exchange-biasing field (H_{ex}) on IrMn thickness (t_{IrMn}) in cases A and B, and under conditions (a) to (c). (b) Dependence of interfacial energy (J_k) on IrMn thickness (t_{IrMn}) in cases A and B.

come the spin rotation barrier, so making spin rotation more facile [17,18].

Fig. 3 plots $H_{\rm c}$ versus $t_{\rm IrMn}$ in the two cases A (without a Ta seed layer) and B (with a Ta seed layer). In general, $H_{\rm c}$ increases as $t_{\rm IrMn}$ increases from 15 Å to 30 Å and then falls as $t_{\rm IrMn}$ increases further, because of the spin coupling and decoupling interaction that

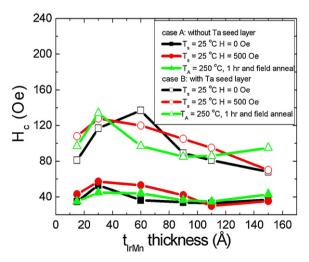
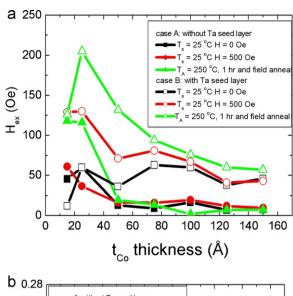


Fig. 3. Coercivity (H_c) versus IrMn thickness (t_{IrMn}) in cases A and B.



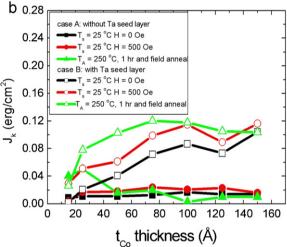


Fig. 4. (a) Dependence of exchange-biasing field $(H_{\rm ex})$ on Co thickness $(t_{\rm Co})$ in cases A and B. (b) Dependence of the interfacial energy $(J_{\rm k})$ on Co thickness $(t_{\rm Co})$ in cases A and B.

occur between the Co and IrMn layers at the Co/IrMn interface as t_{IrMn} is slowly increased [19]. Coupling drag occurs at the Co/IrMn interface, causing H_{ex} to increase as t_{IrMn} is increases from 15 Å to 30 Å. The coupling force between Co and nearby IrMn spins at the interface, exceeds the force between neighboring IrMn spins. The external field (H) must induce both the Co spins and IrMn spins, suggesting that $H_{\rm c}$ increases as $t_{\rm lrMn}$ increases from 15 Å to 30 Å. Furthermore, as t_{IrMn} continues to increase, H_c reaches a maxima, then declines, this is due to decoupling between the interfacial Co spin and the IrMn spin on top. At $t_{\text{IrMn}} \ge 30 \,\text{Å}$, when H_{ex} is fully develop, it indicates that the exchange coupling of interfacial Co and IrMn spins becomes weaker. Therefore, IrMn spin at the interface is strongly pinned by the IrMn spins beneath, H_c switches only the Co spin, and no longer switches the neighboring IrMn spin. However, other researchers proposed an alternative explanation for the increase in H_c that occurs with increasing seed layer thickness [20]. Increasing seed layer thickness produces greater roughness, and this results in exchange-decoupled grain and so increases H_c [20]. Moreover, the H_c of an un-seed layer in the Co/IrMn system shows a similar, albeit smaller trend in H_c .

Fig. 4(a) plots the dependence of $H_{\rm ex}$ on the thickness ($t_{\rm Co}$) of Co in cases A (without a Ta seed layer) and B (with a Ta seed layer). From Eq. (1), $H_{\rm ex}$ (equal to $J_{\rm k}/(M_{\rm s}t_{\rm Co})$) is easily derived. Fig. 4(a) shows that, since ($J_{\rm k}/M_{\rm s}$) is constant, then $H_{\rm ex}$ is proportional to (1/ $t_{\rm Co}$). Therefore, the $H_{\rm ex}$ values with the Ta seed layer are almost

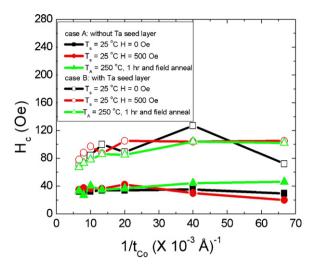


Fig. 5. Coercivity (H_c) versus the inverse of Co thickness $(1/t_{Co})$ in cases A and B, obtained by VSM.

slightly larger than those without a Ta seed-layer are, except for the beginning of two square-data curves. It can be possibly concluded that the random spins arrangement of thinner Co thickness exists in the Ta seed layer, results in lower $H_{\rm ex}$. Finally, from this figure, the Ta-seeded ${\rm Co}(25\,{\rm Å})/{\rm IrMn}(90\,{\rm Å})$ sample has the highest $H_{\rm ex}$ value, approximately 205 Oe. Fig. 4(b) plots $J_{\rm k}$ versus $t_{\rm Co}$. In the general case, $M_{\rm S}$ tends to saturate with sufficient increase in $t_{\rm Co}$. The dependence of $J_{\rm k}$ on $t_{\rm Co}$, seen in Fig. 4(b) is similar to the trend in $M_{\rm S}$ on $t_{\rm Co}$ seen in Fig. 4(a), and similarly, $J_{\rm k}$ tends to saturate slowly as $t_{\rm Co}$ increases. According to calculation, based on Eq. (1), the $J_{\rm k}$ has an optimal value of 0.11 erg/cm² for the post-annealed Ta-seeded ${\rm Co}(75\,{\rm Å})/{\rm IrMn}(90\,{\rm Å})$ sample. In addition, the deposited-field Ta-seeded ${\rm Co}(150\,{\rm Å})/{\rm IrMn}(90\,{\rm Å})$ sample also has an optimal value, since the thicker Co thickness exists a higher saturation magnetization, then $J_{\rm k}$ is enhanced to larger value.

Fig. 5 plots the dependence of $H_{\rm C}$ on $1/t_{\rm CO}$, and shows that $H_{\rm C}$ is inversely proportional to $t_{\rm CO}$. Because of the surface pinning effect present at both the Ta/Co and Co/IrMn interfaces in the Ta-seeded devices, reduction in Co thickness ($t_{\rm CO}$), induces an increase in $H_{\rm C}$, due to an increased contribution to the pinning effect at the Co/IrMn interface from the Ta/Co interface [21]. The surface pinning effect of Ta/Co and Co/IrMn interfaces possibly exists some defects, including oxide, rough roughness, and inclusions. The defects can induce the displacement of domain walls motion difficulty and the $H_{\rm C}$ becomes higher values.

4. Conclusion

The variables H_{ex} , J_{k} and H_{c} of Ta-seeded and un-seeded Co/IrMn were investigated. Together, post-deposition annealing

treatment, and the formation of a Ta seed layer on the Co/lrMn system create the strongest IrMn (111) texture and the highest value for the exchange-bias field, $H_{\rm ex}$. Adding a Ta seed layer alone, does not always result in stronger IrMn (111) texturing, although it can increase $H_{\rm ex}$. Additionally, $H_{\rm c}$ values of Ta-seeded Co/IrMn devices are significantly associated with $t_{\rm IrMn}$, which is explained by coupling and decoupling interactions. The $H_{\rm c}$ values of Ta-seeded and non-seeded devices in the Co/IrMn system are inversely proportional to $t_{\rm Co}$, because of surface pinning in the Ta/Co and Co/IrMn interfaces. Finally, the optimal maximum values for $H_{\rm ex}$ and $J_{\rm k}$ in this investigation are 205 Oe and 0.11 erg/cm², which occur for post-annealed Ta-seeded Co(25 Å)/IrMn(90 Å) and Co(75 Å)/IrMn(90 Å) devices, respectively.

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